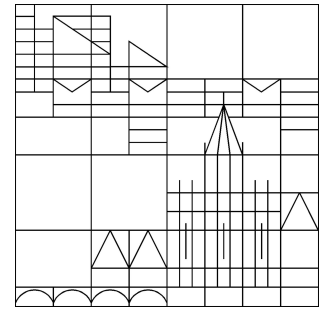


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Asymptotic Behavior of Discontinuous Solutions in 3-D Thermoelasticity with Second Sound

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Abstract: This paper is devoted to the study of the Cauchy problem for linear and semilinear thermoelastic systems with second sound in three space dimensions with discontinuous initial data. Due to Cattaneo's law, replacing Fourier's law for heat conduction, the thermoelastic system with second sound is hyperbolic. We investigate the behavior of discontinuous solutions as the relaxation parameter tends to zero, which corresponds to a formal convergence of the system to the hyperbolic-parabolic type of classical thermoelasticity. By studying expansions with respect to the relaxation parameter of the jumps of the potential part of the system on the evolving characteristic surfaces, we obtain that the jump of the temperature goes to zero while the jumps of the heat flux and the gradient of the potential part of the elastic wave are propagated along the characteristic curves of the elastic fields when the relaxation parameter goes to zero. An interesting phenomenon is when time goes to infinity, the behavior will depend on the mean curvature of the initial surface of discontinuity. These jumps decay exponentially when time goes to infinity, more rapidly for smaller heat conductive coefficient in linear problems and in nonlinear problems when certain growth conditions are imposed on the nonlinear functions.

AMS Subject Classification: 35B40, 74F05

Keywords: hyperbolic thermoelasticity, semilinear, discontinuous solutions, asymptotic behavior, curvature

1 Introduction

Consider the following system of semilinear thermoelasticity with second sound in three space dimensions

$$u_{tt} - \mu \Delta u - (\mu + \lambda) \nabla \operatorname{div} u + \beta \nabla \theta = f(u_t, \theta) \quad (1.1)$$

$$\theta_t + \gamma \operatorname{div} q + \delta \operatorname{div} u_t = g(u_t, \theta) \quad (1.2)$$

$$\tau q_t + q + \kappa \nabla \theta = 0 \quad (1.3)$$

where $u = u(t, x) \in \mathbb{R}^3$, $\theta = \theta(t, x) \in \mathbb{R}$, $q = q(t, x) \in \mathbb{R}^3$ represent the elastic displacement, the temperature and the heat flux, respectively, and are functions of $(t, x) \in (0, \infty) \times \mathbb{R}^3$. The

parameters $\mu, \lambda, \beta, \gamma, \delta, \tau, \kappa$ are constants satisfying

$$\mu, \gamma, \tau, \kappa > 0, 2\mu + 3\lambda > 0, \beta\delta > 0$$

cp. [2, 8, 9]. The given functions f and g are assumed to be smooth. For the differential equations (1.1) - (1.3), we impose the following initial conditions:

$$u(t=0) = u^0, u_t(t=0) = u^1, \theta(t=0) = \theta^0, q(t=0) = q^0 \quad (1.4)$$

We investigate the asymptotic behavior of solutions to (1.1) - (1.4) for which $(\nabla u^0, u^1, \theta^0, q^0)$ may have jumps on an initial surface σ given by

$$\sigma = \{x \in \Omega_0 \subset \mathbb{R}^3 \mid \Phi^0(x) = 0\} \quad (1.5)$$

as the level set of a C^2 -function

$$\Phi^0 : \Omega_0 \subset \mathbb{R}^3 \rightarrow \mathbb{R} \text{ with } |\nabla \Phi^0| \equiv 1 \text{ on } \sigma \quad (1.6)$$

defined on an open set $\Omega_0 \subset \mathbb{R}^3$.

Equation (1.3) is Cattaneo's law for heat conduction and turns, as the so-called relaxation parameter $\tau \rightarrow 0$, into Fourier's law

$$q + \kappa \nabla \theta = 0 \quad (1.7)$$

The system (1.1), (1.2), (1.7) is the hyperbolic-parabolic system of classical thermoelasticity, cp. [6]. The system to be studied here, (1.1) - (1.3), is purely hyperbolic, cp. [2, 8, 9].

The propagation of singularities in the Cauchy problems for the hyperbolic-parabolic coupled system (1.1), (1.2), (1.7) of classical thermoelasticity has been studied in [3, 10, 13, 14] and references therein. It is observed that the propagation of singularities in this coupled system (even more general nonlinear coupled systems) is dominated by the hyperbolic operators. One of main tools used in these works is the microlocal analysis, which requires certain regularity of solutions in general. Thus, the above mentioned works only deal with some weak singularities, in particular, one of important waves, discontinuous solutions do not been investigated.

As pointed out above, formally, when $\tau \rightarrow 0$, the hyperbolic model (1.1), (1.2), (1.3) turns into the hyperbolic-parabolic system. We shall be particularly interested in the behavior of the propagating jumps as $\tau \rightarrow 0$, which should give us some information on the propagation of strong singularities for the hyperbolic-parabolic coupled system. This problem was first studied by authors in [11] for the one-dimensional case. We shall study the multi-dimensional problem in this work.

By decomposing (u, q) into its solenoidal part (u^s, q^s) ($\operatorname{div} u^s = \operatorname{div} q^s = 0$) and its potential part (u^p, q^p) ($\operatorname{rot} u^p = \operatorname{rot} q^p = 0$), the representations and expansions for $\tau \rightarrow 0$ and $t \rightarrow \infty$ of the jumps on the evolving characteristic surfaces are given for the linear problem (Theorem 3.3) and a nonlinear problem (Theorem 4.1). In particular, similar to [11], we shall obtain that the jump of the temperature goes to zero (showing a smoothing effect in the hyperbolic-parabolic coupled systems) while the jumps of the heat flux and the gradient of the potential part of

the elastic wave are propagated along the characteristic curves of the elastic fields when the relaxation parameter τ vanishes.

Moreover, we shall find a very interesting phenomenon whether the jumps of solutions on the hyperbolic characteristic surfaces decay exponentially when $t \rightarrow +\infty$ depends on which one is dominated between the quantity $\frac{\beta\delta}{2\kappa\gamma\sqrt{2\mu+\gamma}}$ and the mean curvature of the initial surface of discontinuity, which is in-visible in the one-dimensional work [11]. Under certain assumptions, these jumps decay exponentially when time goes to infinity, more rapidly for smaller heat conductive coefficient, which is similar to a phenomenon observed by Hoff in [4] for discontinuous solutions to the compressible Navier-Stokes equations.

We remark that our considerations could be extended to variable coefficients and the case that nonlinearities f, q depend on u as well. The case that f, g depending on other quantities like $\nabla\theta$ and ∇u would require further analysis. The two-dimensional case can be studied analogously.

The paper is organized as follows: In Section 2 we give the general setting and a decomposition. The linearized problem ($f = 0, g = 0$) is discussed in Section 3, and Section 4 presents the results on a semilinear problem.

2 Setting and decomposition

Consider the following Cauchy problem

$$u_{tt} - \mu\Delta u - (\mu + \lambda)\nabla \operatorname{div} u + \beta\nabla\theta = f(u_t, \theta) \quad (2.1)$$

$$\theta_t + \gamma \operatorname{div} q + \delta \operatorname{div} u_t = g(u_t, \theta) \quad (2.2)$$

$$\tau q_t + q + \kappa\nabla\theta = 0 \quad (2.3)$$

on $(0, \infty) \times \mathbb{R}^3$,

$$u(t=0) = u^0, u_t(t=0) = u^1, \theta(t=0) = \theta^0, q(t=0) = q^0 \quad (2.4)$$

with data $(\nabla u^0, u^1, \theta^0, q^0)$ being smooth away from and possibly having jumps on a given smooth surface

$$\sigma = \{x \in \Omega_0 \subset \mathbb{R}^3 \mid \Phi^0(x) = 0\}$$

described by a C^2 -function $\Phi^0 : \Omega_0 \subset \mathbb{R}^3 \rightarrow \mathbb{R}$, with Ω_0 open, satisfying

$$|\nabla\Phi^0(x)| = 1 \quad \text{on} \quad \sigma. \quad (2.5)$$

Assumption (2.5) is important for the following techniques used, but is made without loss of generality, since for a surface σ one can choose the distance function

$$\Phi^0(x) := \operatorname{dist}(x, \sigma)$$

which satisfies (2.5), cp. [5, pp. 354-355].

Simple examples are of course spheres of radius $r > 0$ with $\Phi^0(x) = \frac{1}{2r}(|x|^2 - r^2)$ or ellipsoids.

Using the decomposition of vector fields according to the orthogonal decomposition

$$(L^2(\mathbb{R}^3))^3 = \overline{\nabla W^{1,2}(\mathbb{R}^3)} \oplus D_0(\mathbb{R}^3) \quad (2.6)$$

where $W^{1,2}(\mathbb{R}^3)$ denotes the usual Sobolev space, and $D_0(\mathbb{R}^3)$ denotes the set of vector fields having divergence zero, we have

$$u = u^p + u^s, \quad q = q^p + q^s$$

and (2.1) - (2.3) is decomposed into

$$u_{tt}^s - \mu \Delta u^s = f^s(u_t, \theta) \quad (2.7)$$

$$u^s(t=0) = u^{0,s}, \quad u_t^s(t=0) = u^{1,s} \quad (2.8)$$

and

$$\tau q_t^s + q^s = 0 \quad (2.9)$$

$$q^s(t=0) = q^{0,s} \quad (2.10)$$

and

$$u_{tt}^p - \alpha^2 \Delta u^p + \beta \nabla \theta = f^p(u_t, \theta) \quad (2.11)$$

$$\theta_t + \gamma \operatorname{div} q^p + \delta \operatorname{div} u_t^p = g(u_t, \theta) \quad (2.12)$$

$$\tau q_t^p + q^p + \kappa \nabla \theta = 0 \quad (2.13)$$

$$u^p(t=0) = u^{0,p}, \quad u_t^p(t=0) = u^{1,p}, \quad \theta(t=0) = \theta^0, \quad q^p(t=0) = q^{0,p} \quad (2.14)$$

where $\alpha := \sqrt{2\mu + \lambda}$.

Since (2.9), (2.10) is explicitly solvable, and the propagation of singularities for pure wave equations like (2.7), (2.8) is well understood, cp. [1, 7, 12], we concentrate on the fully coupled system (2.11) - (2.14). Let

$$U = (U_1, U_2, U_3, U_4)' := (\operatorname{div} u^p, u_t^p, \theta, q^p)'$$

with the prime denoting the transpose of a vector. Then (2.11) - (2.14) turns into the following system for U :

$$\partial_t U + \sum_{j=1}^3 A_j \partial_j U + A_0 U = F \quad (2.15)$$

$$U(t=0) = U^0 \quad (2.16)$$

where $\partial_t = \frac{\partial}{\partial t}$, $\partial_j = \frac{\partial}{\partial x_j}$, $j = 1, 2, 3$,

$$\sum_{j=1}^3 A_j \partial_j \equiv \begin{pmatrix} 0 & -\operatorname{div} & 0 & 0 \\ -\alpha^2 \nabla & 0 & \beta \nabla & 0 \\ 0 & \delta \operatorname{div} & 0 & \gamma \operatorname{div} \\ 0 & 0 & \frac{\kappa}{\tau} \nabla & 0 \end{pmatrix} \quad (2.17)$$

giving A_j implicitly, and

$$A_0 := \begin{pmatrix} 0_{5 \times 5} & 0_{5 \times 3} \\ 0_{3 \times 5} & \frac{1}{\tau} Id_{\mathbb{R}^3} \end{pmatrix}, \quad F := (0, f^p, g, 0)'. \quad (2.18)$$

3 The linearized problem

We now assume $F = 0$, i.e. we consider the linearized version of (2.15), (2.16). For $\xi \in \mathbb{R}^3 \setminus \{0\}$ we have that the characteristic polynomial

$$\det(\lambda Id_{\mathbb{R}^3} - \sum_{j=1}^3 \xi_j A_j) = \lambda^4 (\lambda^4 - \lambda^2 (\alpha^2 + \beta\delta + \frac{\kappa\gamma}{\tau}) |\xi|^2 + \frac{\kappa\gamma\alpha^2}{\tau} |\xi|^4) \quad (3.1)$$

has real roots $\lambda_k, 1 \leq k \leq 8$, satisfying

$$\lambda_k = 0, \quad 1 \leq k \leq 4 \quad (3.2)$$

$$\lambda_k^2 = \frac{1}{2} (\alpha^2 + \beta\delta + \frac{\kappa\gamma}{\tau} \pm \sqrt{(\alpha^2 + \beta\delta + \frac{\kappa\gamma}{\tau})^2 - \frac{4\kappa\gamma\alpha^2}{\tau}}) |\xi|^2 \equiv \mu_k^2 |\xi|^2, \quad 5 \leq k \leq 8. \quad (3.3)$$

The last factor in (3.1) is similar to the characteristic polynomial in one space dimension, cp. [11], and we know

$$\lambda_k \neq \lambda_j \quad \text{for} \quad 5 \leq k \neq j \leq 8. \quad (3.4)$$

The characteristic surfaces $\Sigma_k = \{(t, x) \mid \Phi_k(t, x) = 0\}, 1 \leq k \leq 8$, evolving from the initial surface $\sigma = \{(0, x) \mid \Phi^0(x) = 0\}$ are determined by

$$\partial_t \Phi_k - \mu_k |\nabla \Phi_k| = 0 \quad (3.5)$$

$$\Phi_k(0, \cdot) = \Phi^0 \quad (3.6)$$

which is solved by

$$\Phi_k(t, x) = \Phi^0(x), \quad 1 \leq k \leq 4 \quad (3.7)$$

$$\Phi_k(t, x) = \mu_k t + \Phi^0(x), \quad 5 \leq k \leq 8 \quad (3.8)$$

where we used the normalization (2.5), $|\nabla \Phi^0| = 1$. Hence

$$\Sigma_k = \{(t, x) \mid \mu_k t + \Phi^0(x) = 0\}. \quad (3.9)$$

Under the assumption (2.5) we have that the eigenvalues of

$$B := \sum_{j=1}^3 \partial_j \Phi^0 A_j \quad (3.10)$$

are

$$\lambda_k = 0, \quad 1 \leq k \leq 4 \quad (3.11)$$

$$\lambda_k = \mu_k = \pm \left\{ \frac{1}{2} \left(\alpha^2 + \beta\delta + \frac{\kappa\gamma}{\tau} \right) \pm \frac{1}{2} \sqrt{\left(\alpha^2 + \beta\delta + \frac{\kappa\gamma}{\tau} \right)^2 - \frac{4\kappa\gamma\alpha^2}{\tau}} \right\}^{\frac{1}{2}}, \quad 5 \leq k \leq 8 \quad (3.12)$$

taking

$$\mu_8 < \mu_6 < 0 < \mu_7 < \mu_5 \quad (3.13)$$

Let $\nabla_{\tau_1}\Phi^0$ and $\nabla_{\tau_2}\Phi^0$ denote two orthogonal vectors on the tangential plane to the initial surface σ . Then the associated right and left eigenvectors of B are the column vectors r_k and the row vectors l_k ,

$$r_1 = c_1(0, (\nabla_{\tau_1}\Phi^0)', 0_{1 \times 4})', \quad r_2 = c_2(0, (\nabla_{\tau_2}\Phi^0)', 0_{1 \times 4})' \quad (3.14)$$

$$r_3 = c_3(0_{1 \times 5}, (\nabla_{\tau_1}\Phi^0)'), \quad r_4 = c_4(0_{1 \times 5}, (\nabla_{\tau_2}\Phi^0)'), \quad (3.15)$$

$$r_k = \left(\frac{\beta}{\alpha^2 - \lambda_k^2}, -\frac{\beta\lambda_k}{\alpha^2 - \lambda_k^2} (\nabla\Phi^0)', 1, \frac{\kappa}{\tau\lambda_k} (\nabla\Phi^0)' \right)', \quad 5 \leq k \leq 8 \quad (3.16)$$

$$l_1 = (0, (\nabla_{\tau_1}\Phi^0)', 0_{1 \times 4}), \quad l_2 = (0, (\nabla_{\tau_2}\Phi^0)', 0_{1 \times 4}) \quad (3.17)$$

$$l_3 = (0_{1 \times 5}, (\nabla_{\tau_1}\Phi^0)'), \quad l_4 = (0_{1 \times 5}, (\nabla_{\tau_2}\Phi^0)'), \quad (3.18)$$

$$l_k = c_k \left(\frac{\alpha^2\delta}{\alpha^2 - \lambda_k^2}, -\frac{\delta\lambda_k}{\alpha^2 - \lambda_k^2} (\nabla\Phi^0)', 1, \frac{\gamma}{\lambda_k} (\nabla\Phi^0)' \right), \quad 5 \leq k \leq 8 \quad (3.19)$$

where $c_k, 1 \leq k \leq 8$, is chosen such that the normalization condition

$$l_j r_k = \begin{cases} 1, & j = k \\ 0, & j \neq k \end{cases} \quad (3.20)$$

holds. For example, we can set

$$c_1 = c_3 = \frac{1}{|\nabla_{\tau_1}\Phi^0|^2}, \quad c_2 = c_4 = \frac{1}{|\nabla_{\tau_2}\Phi^0|^2} \quad (3.21)$$

and

$$1 = c_k \left(\frac{\alpha^2\beta\delta}{(\alpha^2 - \lambda_k^2)^2} + \frac{\beta\delta\lambda_k^2}{(\alpha^2 - \lambda_k^2)^2} + 1 + \frac{\kappa\gamma}{\tau\lambda_k^2} \right), \quad 5 \leq k \leq 8. \quad (3.22)$$

If $k = 5, 6$ we have the expansion, as $\tau \rightarrow 0$,

$$\mu_k^2 = \lambda_k^2 = \frac{1}{2} \left(\alpha^2 + \beta\delta + \frac{\kappa\gamma}{\tau} + \sqrt{\left(\alpha^2 + \beta\delta + \frac{\kappa\gamma}{\tau} \right)^2 - \frac{4\kappa\gamma\alpha^2}{\tau}} \right) = \frac{\kappa\gamma}{\tau} + \beta\delta + \frac{\alpha^2\beta\delta}{\kappa\gamma}\tau + \mathcal{O}(\tau^2) \quad (3.23)$$

(cp. [11]) and for $k = 7, 8$

$$\mu_k^2 = \lambda_k^2 = \frac{1}{2} \left(\alpha^2 + \beta\delta + \frac{\kappa\gamma}{\tau} - \sqrt{\left(\alpha^2 + \beta\delta + \frac{\kappa\gamma}{\tau} \right)^2 - \frac{4\kappa\gamma\alpha^2}{\tau}} \right) = \alpha^2 - \frac{\alpha^2\beta\delta}{\kappa\gamma}\tau + \mathcal{O}(\tau^2). \quad (3.24)$$

Thus for $k = 5, 6$ formula (3.22) reads

$$\begin{aligned} 1 &= c_k \left(1 + \frac{\kappa\gamma}{\kappa\gamma + \beta\delta\tau + \frac{\alpha^2\beta\delta}{\kappa\gamma}\tau^2 + \mathcal{O}(\tau^3)} + \beta\delta \frac{\alpha^2 + \beta\delta + \frac{\kappa\gamma}{\tau} + \frac{\alpha^2\beta\delta}{\kappa\gamma}\tau + \mathcal{O}(\tau^2)}{\left(\frac{\kappa\gamma}{\tau} + \beta\delta - \alpha^2 + \frac{\alpha^2\beta\delta}{\kappa\gamma}\tau + \mathcal{O}(\tau^2) \right)^2} \right) \\ &= c_k(2 + \mathcal{O}(\tau^2)) \end{aligned} \quad (3.25)$$

which implies

$$c_k = \frac{1}{2} + \mathcal{O}(\tau^2), \quad k = 5, 6 \quad (3.26)$$

while for $k = 7, 8$ we get

$$\begin{aligned} 1 &= c_k \left(1 + \frac{\kappa\gamma}{\alpha^2\tau - \frac{\alpha^2\beta\delta}{\kappa\gamma}\tau^2 + \mathcal{O}(\tau^3)} + \beta\delta \frac{2\alpha^2 - \frac{\alpha^2\beta\delta}{\kappa\gamma}\tau + \mathcal{O}(\tau^2)}{\left(\frac{\alpha^2\beta\delta}{\kappa\gamma}\tau + \mathcal{O}(\tau^2)\right)^2} \right) \\ &= c_k \frac{2\kappa^2\gamma^2}{\alpha^2\beta\delta} \tau^{-2} (1 + \mathcal{O}(\tau)) \end{aligned} \quad (3.27)$$

which gives

$$c_k = \frac{\alpha^2\beta\delta}{2\kappa^2\gamma^2} \tau^2 + \mathcal{O}(\tau^3), \quad k = 7, 8. \quad (3.28)$$

Let the matrices L and R be given by

$$L := \begin{pmatrix} l_1 \\ \vdots \\ l_8 \end{pmatrix}, \quad R := (r_1, \dots, r_8).$$

Then $V := LU$, with U satisfying (2.15), (2.16), satisfies

$$\partial_t V + \sum_{j=1}^3 (LA_j R) \partial_j V + \tilde{A}_0 V = \tilde{F} \quad (3.29)$$

$$V(t=0) = V^0 := LU^0 \quad (3.30)$$

where

$$\tilde{A}_0 := LA_0 R + \sum_{j=1}^3 LA_j \partial_j R \quad (3.31)$$

$$\tilde{F} := LF \quad (= 0 \text{ in this section}). \quad (3.32)$$

Then we have as the first result, where $[H]_{\Sigma_k}$ denotes the jump of H along Σ_k :

Lemma 3.1: *Let $V_I := (V_1, V_2, V_3, V_4)$. Then,*

$$[V_I]_{\Sigma_k} = 0, \quad 5 \leq k \leq 8,$$

$$[V_j]_{\Sigma_k} = 0, \quad 5 \leq j \leq 8, \quad k = 1, \dots, 8, \quad j \neq k$$

i.e., V_I are continuous at $\cup_{k=5}^8 \Sigma_k$, and $V_j, j = 5, \dots, 8$, does not have any jump on $\Sigma_k, k = 1, \dots, 8$ for $k \neq j$.

PROOF: Let V be a bounded piecewise smooth solution to (3.29), (3.30) with possible discontinuities only on $\Sigma_k, k = 0, 5, 6, 7, 8$, where $\Sigma_0 := \{(t, x) \in \Omega \mid \Phi^0(x) = 0\}$.

For a fixed $k \in \{5, 6, 7, 8\}$ let Ω be a bounded domain in $\mathbb{R}^+ \times \mathbb{R}^3$ containing only a part of Σ_k while $\Sigma_j \cap \Omega = \emptyset$ for $j \neq k$. Let

$$\Omega^\pm = \{(t, x) \in \Omega \mid \pm \Phi_k(t, x) > 0\}.$$

The equation (3.29) holds in Ω in the distribution sense, that is we have for any function $\varphi \in (C_0^\infty(\Omega))^8$

$$\int_{\Omega} (V \cdot \partial_t \varphi + \sum_{j=1}^3 \partial_j (\varphi^* L A_j R) V) dx dt = - \int_{\Omega} (\tilde{F} - \tilde{A}_0 V) \cdot \varphi dx dt.$$

Using the Gauss integral theorem for the left-hand side on Ω^\pm , respectively, and observing that (3.29) holds in the classical sense in Ω^\pm , it follows

$$\int_{\Omega \cap \Sigma_k} ([V]_{\Sigma_k} \vec{n}_0 + \sum_{j=1}^3 L A_j R [V]_{\Sigma_k} \vec{n}_j) \cdot \varphi = 0$$

where $\vec{n} = (\vec{n}_0, \vec{n}_1, \vec{n}_2, \vec{n}_3)'$ denotes the unit normal vector on Σ_k . Since $\Sigma_k = \{(t, x) \mid \Phi_k(t, x) = 0\}$ we have

$$\vec{n} \parallel (\lambda_k, -(\nabla \Phi^0)').$$

This implies, since φ is arbitrary in $(C_0^\infty(\Omega))^8$,

$$(\lambda_k Id_{\mathbb{R}^8} - \sum_{j=1}^3 L A_j R \partial_j \Phi^0) [V]_{\Sigma_k} = 0$$

or

$$\left(L(\lambda_k Id_{\mathbb{R}^8} - \sum_{j=1}^3 (\partial_j \Phi^0) A_j) R \right) [V]_{\Sigma_k} = 0 \quad (3.33)$$

hence, by definitions of L and R ,

$$(\lambda_k Id_{\mathbb{R}^8} - \tilde{\Lambda}) [V]_{\Sigma_k} = 0 \quad (3.34)$$

where

$$\tilde{\Lambda} := \text{diag} (\lambda_1, \dots, \lambda_8)$$

This implies

$$[V_j]_{\Sigma_k} = 0$$

for $k = 5, \dots, 8$ and $j = 1, \dots, 4$, or $k = 1, \dots, 8$ and $j = 5, \dots, 8$ and $j \neq k$. Q.e.d.

From Lemma 3.1, we know that it remains to study the behavior of V_k on Σ_k . First, we have

Lemma 3.2: *For $1 \leq k \leq 4$ we have*

$$[V_k]_{\Sigma_0} = 0.$$

PROOF: By definition we have

$$V_k = l_k U, \quad k = 1, \dots, 4$$

or

$$V_1 = \partial_t(u^p \cdot \nabla_{\tau_1} \Phi^0), \quad V_2 = \partial_t(u^p \cdot \nabla_{\tau_2} \Phi^0), \quad V_3 = q^p \cdot \nabla_{\tau_1} \Phi^0, \quad V_4 = q^p \cdot \nabla_{\tau_2} \Phi^0 \quad (3.35)$$

For $\omega \in \{u^p, q^p\}$ we have $\text{rot } \omega = 0$, classically away from Σ_0 , and always in the distributional sense, that is for a bounded domain $G \subset \mathbb{R}_t^+ \times \mathbb{R}^3$ with $\Sigma_0 \cap G \neq \emptyset$ and $\varphi \in C_0^\infty(G)$ we have

$$\int_G \omega \text{rot } \varphi = 0$$

implying

$$\int_{\Sigma_0} ([\omega]_{\Sigma_0} \times \vec{\nu}) \varphi = 0$$

where $\vec{\nu}$ denotes the normal on Σ_0 and is parallel to $\nabla \Phi^0$. Hence we get since φ is arbitrary,

$$[w]_{\Sigma_0} \times \vec{\nu} = 0$$

which gives

$$[w]_{\Sigma_0} = 0 \quad \text{or} \quad [w]_{\Sigma_0} \parallel \vec{\nu} \quad (3.36)$$

the latter implying

$$[w \cdot \nabla_{\tau_m} \Phi^0]_{\Sigma_0} = 0, \quad m = 1, 2 \quad (3.37)$$

The assertion of the lemma now follows from (3.35) - (3.37). Q.e.d.

Next, let us derive the evolutional equations of $[V_k]_{\Sigma_k}$ for $5 \leq k \leq 8$.

Denote by D_{mk} the m -th row and k -th column element of a (8×8) -matrix D . Obviously, the k -th equation given in (3.29) can be explicitly written as

$$\partial_t V_k + \sum_{m=1}^8 \sum_{j=1}^3 (L A_j R)_{km} \partial_j V_m = - \sum_{m=1}^8 (\tilde{A}_0)_{km} V_m + \tilde{F}_k \quad (3.38)$$

for $5 \leq k \leq 8$.

From the fact

$$(\lambda_k Id_{\mathbb{R}^8} - \sum_{j=1}^3 L \partial_j \Phi^0 A_j R)_{kk} = (\lambda_k Id_{\mathbb{R}^8} - \tilde{\Lambda})_{kk} = 0 \quad (3.39)$$

with $\tilde{\Lambda} = \text{diag}(\lambda_1, \dots, \lambda_8)$, we know that the operator $\partial_t + \sum_{j=1}^3 (L A_j R)_{kk} \partial_j$ is tangential to $\Sigma_k = \{\lambda_k t + \Phi^0(x) = 0\}$. Similarly, since all entries in line k of $\lambda_k Id_{\mathbb{R}^8} - \sum_{j=1}^3 L \partial_j \Phi^0 A_j R = \lambda_k Id_{\mathbb{R}^8} - \tilde{\Lambda}$ vanish, we deduce that for $m \neq k$, $(0, (L A_1 R)_{km}, (L A_2 R)_{km}, (L A_3 R)_{km})$ is orthogonal to the normal direction $(-\lambda_k, (\nabla \Phi^0)')$ of Σ_k , which implies that $\sum_{j=1}^3 (L A_j R)_{km} \partial_j$ is tangential to Σ_k when $m \neq k$.

Therefore, by using Lemmas 3.1 and 3.2 in (3.38) we obtain that for each $5 \leq k \leq 8$, $[V_k]_{\Sigma_k}$ satisfies the following transport equation:

$$\left(\partial_t + \sum_{j=1}^3 (LA_j R)_{kk} \partial_{x_j} + (\tilde{A}_0)_{kk} \right) [V_k]_{\Sigma_k} = [\tilde{F}_k]_{\Sigma_k} \quad (3.40)$$

with initial conditions

$$[V_k]_{\Sigma_k}(t=0) = [V_k^0]_{\sigma}. \quad (3.41)$$

In order to determine the behavior of $[V_k]_{\Sigma_k}$ from (3.40)(3.41), it is essential to study $(\tilde{A}_0)_{kk}$, where, by (3.31), \tilde{A}_0 is given by

$$\tilde{A}_0 = LA_0 R + \sum_{j=1}^3 LA_j \partial_j R \quad (3.42)$$

Denoting the row l_k and the column r_k by $l_k = (l_{k1}, \dots, l_{k8})$ resp. $r_k = (r_{1k}, \dots, r_{8k})'$ we have, using (2.18), for $k = 5, \dots, 8$

$$(LA_0 R)_{kk} = \frac{1}{\tau} \sum_{j=6}^8 l_{kj} r_{jk}$$

hence, using (3.16), (3.19),

$$(LA_0 R)_{kk} = \frac{c_k \kappa \gamma}{\tau^2 \lambda_k^2}$$

which implies for $k = 5, 6$, using (3.23), (3.26),

$$(LA_0 R)_{kk} = \left(\frac{1}{2} + \mathcal{O}(\tau^2) \right) \frac{\kappa \gamma}{\tau^2 \left(\frac{\kappa \gamma}{\tau} + \beta \delta + \frac{\alpha^2 \beta \delta}{\kappa \gamma} \tau + \mathcal{O}(\tau^2) \right)} = \frac{1}{2\tau} - \frac{\beta \delta}{2\kappa \gamma} + \mathcal{O}(\tau) \quad (3.43)$$

and for $k = 7, 8$ using (3.24), (3.28),

$$(LA_0 R)_{kk} = \left(\frac{\alpha^2 \beta \delta}{2\kappa^2 \gamma^2} \tau^2 + \mathcal{O}(\tau^3) \right) \frac{\kappa \gamma}{\tau^2 \left(\alpha^2 - \frac{\alpha^2 \beta \delta}{\kappa \gamma} \tau + \mathcal{O}(\tau^2) \right)} = \frac{\beta \delta}{2\kappa \gamma} + \mathcal{O}(\tau) \quad (3.44)$$

Now we compute $(\sum_{j=1}^3 LA_j \partial_j R)_{kk}$. Let \vec{e}_j be the j -th standard unit vector having only zero components besides a 1 in the j -th component. Then

$$A_1 = \begin{pmatrix} 0 & -\vec{e}_1' & 0 & 0 \\ -\alpha^2 \vec{e}_1 & 0 & \beta \vec{e}_1 & 0 \\ 0 & \delta \vec{e}_1' & 0 & \gamma \vec{e}_1' \\ 0 & 0 & \frac{\kappa}{\tau} \vec{e}_1 & 0 \end{pmatrix}$$

and A_2, A_3 arise from A_1 by replacing \vec{e}_1 by \vec{e}_2 and \vec{e}_3 , respectively.

Denote by $LA_j \partial_j R = (a_{ik}^j)_{8 \times 8}$ for any $j = 1, 2, 3$. Then, by a direct computation, we deduce that for $k = 5, \dots, 8$,

$$\begin{aligned} a_{kk}^1 &= -l_{k1} \partial_1 r_{2k} + l_{k2} (\beta \partial_1 r_{5k} - \alpha^2 \partial_1 r_{1k}) + l_{k5} (\delta \partial_1 r_{2k} + \gamma \partial_1 r_{6k}) + l_{k6} \frac{\kappa}{\tau} \partial_1 r_{5k} \\ a_{kk}^2 &= -l_{k1} \partial_2 r_{3k} + l_{k3} (\beta \partial_2 r_{5k} - \alpha^2 \partial_2 r_{1k}) + l_{k5} (\delta \partial_2 r_{3k} + \gamma \partial_2 r_{7k}) + l_{k7} \frac{\kappa}{\tau} \partial_2 r_{5k} \\ a_{kk}^3 &= -l_{k1} \partial_3 r_{4k} + l_{k4} (\beta \partial_3 r_{5k} - \alpha^2 \partial_3 r_{1k}) + l_{k5} (\delta \partial_3 r_{4k} + \gamma \partial_3 r_{8k}) + l_{k8} \frac{\kappa}{\tau} \partial_3 r_{5k} \end{aligned}$$

which gives

$$\begin{aligned} \left(\sum_{j=1}^3 LA_j \partial_j R \right)_{kk} &= -l_{k1}(\partial_1 r_{2k} + \partial_2 r_{3k} + \partial_3 r_{4k}) \\ &\quad + l_{k5} \left(\delta(\partial_1 r_{2k} + \partial_2 r_{3k} + \partial_3 r_{4k}) + \gamma(\partial_1 r_{6k} + \partial_2 r_{7k} + \partial_3 r_{8k}) \right). \end{aligned}$$

Using (3.16) again we conclude

$$\left(\sum_{j=1}^3 LA_j \partial_j R \right)_{kk} = \left((-\delta l_{k5} + l_{k1}) \frac{\beta \lambda_k}{\alpha^2 - \lambda_k^2} + l_{k5} \frac{\kappa \gamma}{\tau \lambda_k} \right) \Delta \Phi^0 \quad (3.45)$$

where Δ denotes the Laplace operator in \mathbb{R}^3 . Since

$$\delta l_{k5} - l_{k1} = c_k \left(\delta - \frac{\alpha^2 \delta}{\alpha^2 - \lambda_k^2} \right) = -c_k \frac{\delta \lambda_k^2}{\alpha^2 - \lambda_k^2}$$

we get from (3.45)

$$\left(\sum_{j=1}^3 LA_j \partial_j R \right)_{kk} = c_k \left(\frac{\beta \delta \lambda_k^3}{(\alpha^2 - \lambda_k^2)^2} + \frac{\kappa \gamma}{\tau \lambda_k} \right) \Delta \Phi^0. \quad (3.46)$$

Using the expansions (3.23), (3.26) we obtain for $k = 5, 6$

$$\begin{aligned} c_k \left(\frac{\beta \delta \lambda_k^3}{(\alpha^2 - \lambda_k^2)^2} + \frac{\kappa \gamma}{\tau \lambda_k} \right) &= \left(\frac{1}{2} + \mathcal{O}(\tau^2) \right) \frac{1}{\lambda_k} \left(\frac{\kappa \gamma}{\tau} + \beta \delta + \mathcal{O}(\tau) \right) \\ \frac{1}{\lambda_k} &= \pm \sqrt{\frac{\tau}{\kappa \gamma}} (1 + \mathcal{O}(\tau)) \end{aligned}$$

implying

$$\left(\sum_{j=1}^3 LA_j \partial_j R \right)_{kk} = \pm \frac{1}{2} \sqrt{\frac{\kappa \gamma}{\tau}} (1 + \mathcal{O}(\tau)) \Delta \Phi^0, \quad k = 5, 6. \quad (3.47)$$

Similarly we obtain, using the expansions (3.24), (3.28),

$$\left(\sum_{j=1}^3 LA_j \partial_j R \right)_{kk} = \pm \left(\frac{\alpha}{2} + \mathcal{O}(\tau) \right) \Delta \Phi^0, \quad k = 7, 8 \quad (3.48)$$

Combining (3.42), (3.43), (3.44), (3.47) with (3.48), we obtain

$$(\tilde{A}_0)_{kk} = \begin{cases} \frac{1}{2\tau} \pm \frac{1}{2} \sqrt{\frac{\kappa \gamma}{\tau}} \Delta \Phi^0 + \mathcal{O}(1), & k = 5, 6 \\ \frac{\beta \delta}{2\kappa \gamma} \pm \frac{\alpha}{2} \Delta \Phi^0 + \mathcal{O}(\tau), & k = 7, 8 \end{cases} \quad (3.49)$$

From (3.49) we already notice the phenomenon that the mean curvature H of the initial surface σ ,

$$H = \frac{\Delta \Phi^0}{2}$$

will play an essential role in the behavior of the jumps as $t \rightarrow \infty$ or as $\tau \rightarrow 0$.

For any fixed $x^0 \in \sigma$, denote by

$$t \rightarrow (t, x_1(t; 0, x^0), x_2(t; 0, x^0), x_3(t; 0, x^0))$$

the characteristic line of the operator $\partial_t + \sum_{j=1}^3 (LA_j R)_{kk} \partial_{x_j}$ passing through $(0, x^0)$, i.e. $x_j(t; 0, x^0)$ satisfies

$$\begin{cases} \frac{\partial x_j(t; 0, x^0)}{\partial t} = (LA_j R)_{kk}(t, x_1(t; 0, x^0), x_2(t; 0, x^0), x_3(t; 0, x^0)) \\ x_j(0; 0, x^0) = x_j^0, \quad 1 \leq j \leq 3 \end{cases}$$

and

$$[V_k]_{\Sigma_k(t)} = [V_k]_{\Sigma_k}(t, x_1(t; 0, x^0), x_2(t; 0, x^0), x_3(t; 0, x^0)). \quad (3.50)$$

Then, from (3.40) and (3.49) we conclude:

$$\begin{aligned} [V_k]_{\Sigma_k(t)} &= [V_k^0]_{\sigma} e^{-\int_0^t (\tilde{A}_0)_{kk}(x(s; 0, x^0)) ds} \\ &= \begin{cases} e^{-\frac{t}{2\tau} \mp \frac{1}{2} \sqrt{\frac{\kappa\gamma}{\tau}} \int_0^t (\Delta\Phi^0(x(s; 0, x^0)) + \mathcal{O}(\sqrt{\tau})) ds} [V_k^0]_{\sigma}, & k = 5, 6 \\ e^{-\frac{\beta\delta}{2\kappa\gamma} t \mp \frac{\alpha}{2} \int_0^t (\Delta\Phi^0(x(s; 0, x^0)) + \mathcal{O}(\tau)) ds} [V_k^0]_{\sigma}, & k = 7, 8. \end{cases} \end{aligned} \quad (3.51)$$

For $\tau \rightarrow 0$ the dominating term for $k = 5, 6$ is $e^{-\frac{t}{2\tau}}$, i.e. we have exponential decay of the jumps of V_k on Σ_k as $\tau \rightarrow 0$ or $t \rightarrow \infty$ for a fixed small $\tau > 0$. If $k = 7, 8$ the dominating term, for $\tau \rightarrow 0$, is $\exp(-\int_0^t (\frac{\beta\delta}{2\kappa\gamma} \pm \frac{\alpha}{2} \Delta\Phi^0(x(s; 0, x^0))) ds)$, whether the jumps of V_k on Σ_k decay exponentially depends on the size of the mean curvature ($= \Delta\Phi^0/2$).

EXAMPLE 3.1: Let σ be the sphere of radius r :

$$\sigma = \{x \in \mathbb{R}^3 \mid |x| = r\} = \{x \mid \Phi^0(x) \equiv |x| - r = 0\}.$$

Then, we have

$$\Sigma_k = \{(t, x) \mid \mu_k t = r - |x|\} = \{(t, x) \mid |x| = r - \mu_k t\}$$

Remember the convention on the signs of μ_k given in (3.13). Spreading as $t \rightarrow \infty$ are Σ_6, Σ_8 , and

$$\Delta\Phi^0(x_0) = \frac{2}{|x_0|} = \frac{2}{r} > 0.$$

Thus, as $t \rightarrow +\infty$, $[V_6]_{\Sigma_6}$ is decaying exponentially, while $[V_8]_{\Sigma_8}$ decays (grows resp.) exponentially if

$$\frac{\beta\delta}{\alpha\kappa\gamma} > \Delta\Phi^0 = \frac{2}{r} \quad \left(\frac{\beta\delta}{\alpha\kappa\gamma} < \Delta\Phi^0 = \frac{2}{r} \text{ resp. } \right), \quad (3.52)$$

that is depending on the size of the mean curvature $H = \frac{1}{r}$.

To complete the discussion on the linear problem, we have to compute $U = RV$ from V ,

$$U = (\operatorname{div} u^p, u_t^p, \theta, q^p)' = RV = \sum_{k=1}^8 V_k r_k.$$

Using (3.14) - (3.16) we conclude

$$\operatorname{div} u^p = \sum_{k=5}^8 \frac{\beta}{\alpha^2 - \lambda_k^2} V_k \quad (3.53)$$

$$u_t^p = - \sum_{k=5}^8 \frac{\beta \lambda_k}{\alpha^2 - \lambda_k^2} \nabla \Phi^0 V_k + \frac{\nabla_{\tau_1} \Phi^0}{|\nabla_{\tau_1} \Phi^0|^2} V_1 + \frac{\nabla_{\tau_2} \Phi^0}{|\nabla_{\tau_2} \Phi^0|^2} V_2 \quad (3.54)$$

$$\theta = \sum_{k=5}^8 V_k \quad (3.55)$$

$$q^p = \sum_{k=5}^8 \frac{\kappa}{\tau \lambda_k} \nabla \Phi^0 V_k + \frac{\nabla_{\tau_1} \Phi^0}{|\nabla_{\tau_1} \Phi^0|^2} V_3 + \frac{\nabla_{\tau_2} \Phi^0}{|\nabla_{\tau_2} \Phi^0|^2} V_4. \quad (3.56)$$

Observing $[V_j]_{\Sigma_k} = 0$ for $j = 1, \dots, 4$ and $k = 1, \dots, 8$, cp. Lemmas 3.1 and 3.2, we get from (3.53) - (3.56) for $k = 5, \dots, 8$:

$$[\operatorname{div} u^p]_{\Sigma_k} = \frac{\beta}{\alpha^2 - \lambda_k^2} [V_k]_{\Sigma_k} \quad (3.57)$$

$$[u_t^p]_{\Sigma_k} = - \frac{\beta \lambda_k}{\alpha^2 - \lambda_k^2} \nabla \Phi^0 [V_k]_{\Sigma_k} \quad (3.58)$$

$$[\theta]_{\Sigma_k} = [V_k]_{\Sigma_k} \quad (3.59)$$

$$[q^p]_{\Sigma_k} = \frac{\kappa}{\tau \lambda_k} \nabla \Phi^0 [V_k]_{\Sigma_k}. \quad (3.60)$$

Using the expansions

$$\frac{1}{\alpha^2 - \lambda_k^2} = \begin{cases} -\frac{\tau}{\kappa \gamma} + \mathcal{O}(\tau^2), & k = 5, 6 \\ \frac{\kappa \gamma}{\alpha^2 \beta \delta} \frac{1}{\tau} + \mathcal{O}(1), & k = 7, 8 \end{cases} \quad (3.61)$$

$$\lambda_k = \begin{cases} \pm \sqrt{\frac{\kappa \gamma}{\tau}} (1 + \mathcal{O}(\tau)), & k = 5, 6 \\ \pm \alpha + \mathcal{O}(\tau), & k = 7, 8 \end{cases} \quad (3.62)$$

we conclude from (3.57) - (3.60), using (3.51),

$$[\operatorname{div} u^p]_{\Sigma_k} = \begin{cases} -\frac{\beta}{\kappa \gamma} (\tau + \mathcal{O}(\tau^2)) e^{-\frac{t}{2\tau} (1 + \mathcal{O}(\sqrt{\tau}))} [V_k^0]_{\sigma}, & k = 5, 6 \\ \frac{\kappa \gamma}{\alpha^2 \beta \delta} \left(\frac{1}{\tau} + \mathcal{O}(1) \right) e^{-\int_0^t (\frac{\beta \delta}{2\kappa \gamma} \pm \frac{\alpha}{2} \Delta \Phi^0(x(s; 0, x^0)) + \mathcal{O}(\tau)) ds} [V_k^0]_{\sigma}, & k = 7, 8 \end{cases} \quad (3.63)$$

$$[u_t^p]_{\Sigma_k} = \begin{cases} \pm \frac{\beta}{\sqrt{\kappa \gamma}} \nabla \Phi^0 \sqrt{\tau} (1 + \mathcal{O}(\tau)) e^{-\frac{t}{2\tau} (1 + \mathcal{O}(\sqrt{\tau}))} [V_k^0]_{\sigma}, & k = 5, 6 \\ \mp \frac{\kappa \gamma}{\alpha \delta} \left(\frac{1}{\tau} + \mathcal{O}(1) \right) e^{-\int_0^t (\frac{\beta \delta}{2\kappa \gamma} \pm \frac{\alpha}{2} \Delta \Phi^0(x(s; 0, x^0)) + \mathcal{O}(\tau)) ds} [V_k^0]_{\sigma}, & k = 7, 8 \end{cases} \quad (3.64)$$

$$[\theta]_{\Sigma_k} = \begin{cases} e^{-\frac{t}{2\tau} (1 + \mathcal{O}(\sqrt{\tau}))} [V_k^0]_{\sigma}, & k = 5, 6 \\ e^{-\int_0^t (\frac{\beta \delta}{2\kappa \gamma} \pm \frac{\alpha}{2} \Delta \Phi^0(x(s; 0, x^0)) + \mathcal{O}(\tau)) ds} [V_k^0]_{\sigma}, & k = 7, 8 \end{cases} \quad (3.65)$$

$$[q^p]_{\Sigma_k} = \begin{cases} \pm \sqrt{\frac{\kappa}{\gamma}} \nabla \Phi^0 \left(\frac{1}{\sqrt{\tau}} + \mathcal{O}(\sqrt{\tau}) \right) e^{-\frac{t}{2\tau} (1 + \mathcal{O}(\sqrt{\tau}))} [V_k^0]_{\sigma}, & k = 5, 6 \\ \pm \frac{\kappa}{\alpha} \nabla \Phi^0 \left(\frac{1}{\tau} + \mathcal{O}(1) \right) e^{-\int_0^t (\frac{\beta \delta}{2\kappa \gamma} \pm \frac{\alpha}{2} \Delta \Phi^0(x(s; 0, x^0)) + \mathcal{O}(\tau)) ds} [V_k^0]_{\sigma}, & k = 7, 8 \end{cases} \quad (3.66)$$

We see from (3.63) - (3.66) that if $[V_k^0]_{\sigma} \neq 0$, $k = 5, \dots, 8$ then the jumps of $(\operatorname{div} u^p, u_t^p, \theta, q^p)$ on Σ_k ($k = 5, 6$) decay exponentially as $\tau \rightarrow 0$, and for fixed small τ , they decay exponentially as

well when $t \rightarrow \infty$, while the behavior of those on Σ_k ($k = 7, 8$) depends on the mean curvature ($\frac{1}{2}\Delta\Phi^0$) of the initial surface σ , even allowing exponential growth.

Obviously, from the transformation $V = L \cdot (\operatorname{div}u^p, u_t^p, \theta, q^p)'$, we have

$$V_k = l_{k1}\operatorname{div}u^p + \sum_{j=2}^4 l_{kj}\partial_t u_{j-1}^p + l_{k5}\theta + \sum_{j=6}^8 l_{kj}q_{j-5}^p, \quad 1 \leq k \leq 8$$

which gives

$$\begin{aligned} V_k &= \left(\frac{1}{2} + \mathcal{O}(\tau^2)\right) \left[-\frac{\alpha^2\delta\tau}{\kappa\gamma}(1 + \mathcal{O}(\tau))\operatorname{div}u^p \pm \delta\sqrt{\frac{\tau}{\kappa\gamma}}(1 + \mathcal{O}(\tau))\nabla\Phi^0 \cdot \partial_t u^p\right. \\ &\quad \left.+ \theta \pm \sqrt{\frac{\tau\gamma}{\kappa}}(1 + \mathcal{O}(\tau))\nabla\Phi^0 \cdot q^p\right] \end{aligned} \quad (3.67)$$

for $k = 5, 6$, and

$$\begin{aligned} V_k &= \left(\frac{\alpha^2\beta\delta}{2\kappa^2\gamma^2}\tau^2 + \mathcal{O}(\tau^3)\right) \left[\frac{\kappa\gamma}{\beta\tau}(1 + \mathcal{O}(\tau))\operatorname{div}u^p \mp \frac{\kappa\gamma}{\alpha\beta\tau}(1 + \mathcal{O}(\tau))\nabla\Phi^0 \cdot \partial_t u^p\right. \\ &\quad \left.+ \theta \pm \frac{\gamma}{\alpha}(1 + \mathcal{O}(\tau))\nabla\Phi^0 \cdot q^p\right] \end{aligned} \quad (3.68)$$

for $k = 7, 8$.

Suppose that

$$[\operatorname{div}u^{0,p}]_\sigma, \quad [u^{1,p}]_\sigma, \quad [\theta^0]_\sigma, \quad [q^{0,p}]_\sigma$$

may not vanish for the initial data given in (1.4), then from (3.67),(3.68) we have

$$[V_k^0]_\sigma = \begin{cases} -\frac{\alpha^2\delta\tau}{2\kappa\gamma}(1 + \mathcal{O}(\tau))[\operatorname{div}u^{0,p}]_\sigma \pm \frac{\delta}{2}\sqrt{\frac{\tau}{\kappa\gamma}}(1 + \mathcal{O}(\tau))\nabla\Phi^0 \cdot [u^{1,p}]_\sigma \\ \quad + \left(\frac{1}{2} + \mathcal{O}(\tau)\right)[\theta^0]_\sigma \pm \frac{1}{2}\sqrt{\frac{\tau\gamma}{\kappa}}(1 + \mathcal{O}(\tau))\nabla\Phi^0 \cdot [q^{0,p}]_\sigma, & k = 5, 6 \\ \frac{\alpha^2\delta\tau}{2\kappa\gamma}(1 + \mathcal{O}(\tau))[\operatorname{div}u^{0,p}]_\sigma \mp \frac{\alpha\delta\tau}{2\kappa\gamma}(1 + \mathcal{O}(\tau))\nabla\Phi^0 \cdot [u^{1,p}]_\sigma \\ \quad + \frac{\alpha^2\beta\delta}{2\kappa^2\gamma^2}\tau^2(1 + \mathcal{O}(\tau))[\theta^0]_\sigma \pm \frac{\alpha\beta\gamma}{2\kappa^2\gamma}\tau^2(1 + \mathcal{O}(\tau))\nabla\Phi^0 \cdot [q^{0,p}]_\sigma, & k = 7, 8. \end{cases} \quad (3.69)$$

Substituting (3.69) into (3.63)-(3.66), it follows

$$[\operatorname{div}u^p]_{\Sigma_k} = \begin{cases} \mathcal{O}(\tau)e^{-\frac{t}{2\tau}(1+\mathcal{O}(\sqrt{\tau}))}, & k = 5, 6 \\ \frac{1}{2\alpha} \left(\alpha[\operatorname{div}u^{0,p}]_\sigma \mp \nabla\Phi^0 \cdot [u^{1,p}]_\sigma + \mathcal{O}(\tau) \right) \cdot \\ \quad \cdot e^{-\int_0^t \left(\frac{\beta\delta}{2\kappa\gamma} \pm \frac{\alpha}{2}\Delta\Phi^0(x(s;0,x^0)) + \mathcal{O}(\tau) \right) ds}, & k = 7, 8 \end{cases} \quad (3.70)$$

$$[u_t^p]_{\Sigma_k} = \begin{cases} \mathcal{O}(\sqrt{\tau})e^{-\frac{t}{2\tau}(1+\mathcal{O}(\sqrt{\tau}))}, & k = 5, 6 \\ \mp \frac{1}{2} \left(\alpha[\operatorname{div}u^{0,p}]_\sigma \mp \nabla\Phi^0 \cdot [u^{1,p}]_\sigma + \mathcal{O}(\tau) \right) \cdot \\ \quad \cdot e^{-\int_0^t \left(\frac{\beta\delta}{2\kappa\gamma} \pm \frac{\alpha}{2}\Delta\Phi^0(x(s;0,x^0)) + \mathcal{O}(\tau) \right) ds}, & k = 7, 8 \end{cases} \quad (3.71)$$

$$[\theta]_{\Sigma_k} = \begin{cases} \frac{1}{2}([\theta^0]_\sigma + \mathcal{O}(\sqrt{\tau}))e^{-\frac{t}{2\tau}(1+\mathcal{O}(\sqrt{\tau}))}, & k = 5, 6 \\ \mathcal{O}(\tau)e^{-\int_0^t \left(\frac{\beta\delta}{2\kappa\gamma} \pm \frac{\alpha}{2}\Delta\Phi^0(x(s;0,x^0)) + \mathcal{O}(\tau) \right) ds}, & k = 7, 8 \end{cases} \quad (3.72)$$

$$[q^p]_{\Sigma_k} = \begin{cases} \left(\pm \frac{1}{2}\sqrt{\frac{\kappa}{\gamma\tau}}[\theta^0]_\sigma + \frac{1}{2}\nabla\Phi^0 \cdot \left([q^{0,p}]_\sigma + \frac{\delta}{\gamma}u^{1,p}]_\sigma \cdot \nabla\Phi^0 \right) + \mathcal{O}(\sqrt{\tau}) \right) e^{-\frac{t}{2\tau}(1+\mathcal{O}(\sqrt{\tau}))}, & k = 5, 6 \\ \pm \frac{\delta}{2\gamma} \left((\alpha[\operatorname{div}u^{0,p}]_\sigma \mp \nabla\Phi^0 \cdot [u^{1,p}]_\sigma) \nabla\Phi^0 + \mathcal{O}(\tau) \right) \cdot \\ \quad \cdot e^{-\int_0^t \left(\frac{\beta\delta}{2\kappa\gamma} \pm \frac{\alpha}{2}\Delta\Phi^0(x(s;0,x^0)) + \mathcal{O}(\tau) \right) ds}, & k = 7, 8. \end{cases} \quad (3.73)$$

Altogether the above discussion we conclude

Theorem 3.3: *Suppose that the initial data of $\operatorname{div}u^p, \partial_t u^p, \theta$ and q^p may have jump on $\sigma = \{\Phi^0(x) = 0\}$ with $|\nabla\Phi^0(x)| = 1$, then the propagation of strong singularities of solutions to the linearized problem (2.11)-(2.14) ($f^p = g = 0$) is described by (3.70)-(3.73). In particular, we have*

(1) *The jumps of $\operatorname{div}u^p, \partial_t u^p, \theta, q^p$ on Σ_5 and Σ_6 decay exponentially both when $\tau \rightarrow 0$ for a fixed $t > 0$ and when $t \rightarrow +\infty$ for a fixed $\tau > 0$.*

(2) *The jumps of $\operatorname{div}u^p, \partial_t u^p, q^p$ on Σ_7 (Σ_8 resp.) are propagated, and when $t \rightarrow +\infty$ they will decay exponentially as soon as $\frac{\beta\delta}{\kappa\gamma} + \alpha\Delta\Phi^0$ ($\frac{\beta\delta}{\kappa\gamma} - \alpha\Delta\Phi^0$ resp.) being positive, more rapidly for smaller heat conductive coefficient $\kappa\gamma$, while the jump of the temperature θ on Σ_7 and Σ_8 vanishes of order $O(\tau)$ when $\tau \rightarrow 0$, which shows a smoothing effect in the system (2.11)-(2.13) when the thermoelastic model with second sound converges to the hyperbolic-parabolic type of classical thermoelasticity.*

4 A semilinear problem

In this section, we shall mainly study the asymptotic behavior of discontinuities in the following semilinear problem:

$$u_{tt} - \alpha^2 \Delta u + \beta \nabla \theta = f(u_t, \theta) \quad (4.1)$$

$$\theta_t + \gamma \operatorname{div} q + \delta \operatorname{div} u_t = g(u_t, \theta) \quad (4.2)$$

$$\tau q_t + q + \kappa \nabla \theta = 0 \quad (4.3)$$

on $(0, \infty) \times \mathbb{R}^3$,

$$u(t=0) = u^0, \quad u_t(t=0) = u^1, \quad \theta(t=0) = \theta^0, \quad q(t=0) = q^0 \quad (4.4)$$

with all coefficients being the same as in (2.11)-(2.13), and

$$\operatorname{rot} u = 0, \quad \operatorname{rot} q^0 = 0. \quad (4.5)$$

By comparing the above problem with the problem (2.1)-(2.4), it is easy to see that if we assume that the nonlinear function f in (2.1) is a conserved field, i.e. there is a scalar function \tilde{g} such that $f = \nabla \tilde{g}$, then from the problems (2.7)(2.8) and (2.9)(2.10), we know that the solenoidal part (u^s, q^s) vanishes when the initial data satisfy $\operatorname{rot} u^0 = \operatorname{rot} u^1 = 0$. Therefore, we get that the unknowns (u, θ, q) to the original thermoelastic systems (2.1)-(2.3) of second sound satisfy the above problem. In the remainder, we shall only focus on the problem (4.1)-(4.5).

Denote by $U = (\operatorname{div}u, u_t, \theta, q)'$, and $L = (l'_1, \dots, l'_8)'$ the matrix composed by the left eigenvectors $\{l_j\}_{j=1}^8$ given in (3.17)-(3.19). As in §3, we know that $V = LU$ satisfies the problem (3.29), (3.30):

$$\partial_t V + \sum_{j=1}^3 (LA_j R) \partial_j V + \tilde{A}_0 V = \tilde{F} \quad (4.6)$$

$$V(t=0) = V^0 := LU^0 \quad (4.7)$$

where $\tilde{F}(V) = L \cdot (0, f(V), g(V), 0)'$, with

$$f(V) = f\left(\sum_{k=1}^8 r_{2k} V_k, \dots, \sum_{k=1}^8 r_{5k} V_k\right)$$

and $g(V)$ having the same form.

Obviously, the results stated in Lemma 3.1 still hold for the problem (4.6), (4.7). Thus, $[V_j]_{\Sigma_k} = 0$ for $1 \leq j \leq 4, 5 \leq k \leq 8$, and $5 \leq j \leq 8, 1 \leq k \leq 8$ with $j \neq k$, and the jump of V_k on Σ_k satisfies the transport equation:

$$\left(\partial_t + \sum_{j=1}^3 (LA_j R)_{kk} \partial_{x_j} + (\tilde{A}_0)_{kk}\right) [V_k]_{\Sigma_k} = [\tilde{F}_k]_{\Sigma_k} \quad (4.8)$$

where

$$\tilde{F}_k(V) = \sum_{j=2}^4 l_{kj} f_j(V) + l_{k5} g(V).$$

Suppose that f and g are globally Lipschitz in their arguments. From (3.16), (3.19) we have

$$\begin{aligned} |[\tilde{F}_k]_{\Sigma_k}| &\leq \sum_{l=1}^4 \left\{ \sum_{j=1}^3 |l_{k(j+1)}| \cdot \|\partial_l f_j\|_{L^\infty} + |l_{k5}| \cdot \|\partial_l g\|_{L^\infty} \right\} |r_{(l+1)k}| \cdot |[V_k]_{\Sigma_k}| \\ &\leq |c_k| \left\{ \sum_{l=1}^3 \left| \frac{\beta \lambda_k}{\alpha^2 - \lambda_k^2} \partial_{x_l} \Phi^0 \right| \left(\sum_{j=1}^3 \left| \frac{\delta \lambda_k}{\alpha^2 - \lambda_k^2} \partial_{x_j} \Phi^0 \right| \cdot \|\partial_l f_j\|_{L^\infty} + \|\partial_l g\|_{L^\infty} \right) \right. \\ &\quad \left. + \sum_{j=1}^3 \left| \frac{\delta \lambda_k}{\alpha^2 - \lambda_k^2} \partial_{x_j} \Phi^0 \right| \cdot \|\partial_4 f_j\|_{L^\infty} + \|\partial_4 g\|_{L^\infty} \right\} |[V_k]_{\Sigma_k}| \end{aligned}$$

By using (3.61), (3.62), we have

$$|[\tilde{F}_k(V)]_{\Sigma_k}| \leq \begin{cases} \mathcal{O}(1) |[V_k]_{\Sigma_k}|, & k = 5, 6 \\ \left(\frac{1}{2} \sum_{j,l=1}^3 |\partial_{x_j} \Phi^0 \partial_{x_l} \Phi^0| \|\partial_l f_j\|_{L^\infty} + \mathcal{O}(\tau)\right) |[V_k]_{\Sigma_k}|, & k = 7, 8. \end{cases} \quad (4.9)$$

As in (3.50), for any fixed $x^0 \in \sigma$, denote by

$$t \rightarrow (t, x_1(t; 0, x^0), x_2(t; 0, x^0), x_3(t; 0, x^0))$$

the characteristic line of the operator $\partial_t + \sum_{j=1}^3 (LA_j R)_{kk} \partial_{x_j}$ passing through $(0, x^0)$, and

$$[V_k]_{\Sigma_k(t)} = [V_k]_{\Sigma_k}(t, x_1(t; 0, x^0), x_2(t; 0, x^0), x_3(t; 0, x^0)).$$

Then, from (4.8) we have

$$\partial_t [V_k]_{\Sigma_k(t)} + (\tilde{A}_0)_{kk}(x(t; 0, x^0)) [V_k]_{\Sigma_k(t)} = [\tilde{F}_k]_{\Sigma_k(t)} \quad (4.10)$$

From (4.10) we immediately obtain

$$[V_k]_{\Sigma_k(t)} = [V_k]_{\Sigma_k(0)} e^{-\int_0^t (\tilde{A}_0)_{kk}(x(s; 0, x^0)) ds} + \int_0^t [\tilde{F}_k]_{\Sigma_k(s)} e^{-\int_s^t (\tilde{A}_0)_{kk}(x(s_1; 0, x^0)) ds_1} ds. \quad (4.11)$$

From (4.11) and (4.9), we see that for $k = 5, 6$, the behavior of $[V_k]_{\Sigma_k}$ is the same as given in (3.51).

When $k = 7, 8$, from (4.11) and (4.9), we have

$$|[V_k]_{\Sigma_k(t)}| \leq |[V_k]_{\Sigma_k(0)}| \exp \left(- \int_0^t (\tilde{A}_0)_{kk}(x(s; 0, x^0)) ds + \left(\frac{1}{2} \sum_{j,l=1}^3 \|\partial_l f_j\|_{L^\infty} \|\partial_{x_j} \Phi^0 \partial_{x_l} \Phi^0\|_{L^\infty} + \mathcal{O}(\tau) \right) t \right) \quad (4.12)$$

and

$$\begin{aligned} |[V_k]_{\Sigma_k(t)}| &\geq |[V_k]_{\Sigma_k(0)}| \exp \left(- \int_0^t (\tilde{A}_0)_{kk}(x(s; 0, x^0)) ds \right) - \int_0^t \left(\frac{1}{2} \sum_{j,l=1}^3 \|\partial_l f_j\|_{L^\infty} \|\partial_{x_j} \Phi^0 \partial_{x_l} \Phi^0\|_{L^\infty} \right. \\ &\quad \left. + \mathcal{O}(\tau) \right) |[V_k]_{\Sigma_k(s)}| e^{-\int_s^t (\tilde{A}_0)_{kk}(x(s_1; 0, x^0)) ds_1} ds \end{aligned} \quad (4.13)$$

By using the estimate (4.12) for $[V_k]_{\Sigma_k(s)}$ in the second term on the right hand of (4.13), one gets

$$|[V_k]_{\Sigma_k(t)}| \geq |[V_k]_{\Sigma_k(0)}| e^{-\int_0^t (\tilde{A}_0)_{kk}(x(s; 0, x^0)) ds} \left(2 - e^{\left(\frac{1}{2} \sum_{j,l=1}^3 \|\partial_l f_j\|_{L^\infty} \|\partial_{x_j} \Phi^0 \partial_{x_l} \Phi^0\|_{L^\infty} + \mathcal{O}(\tau) \right) t} \right) \quad (4.14)$$

To obtain the information of $U = (\operatorname{div} u, u_t, \theta, q)'$, we follow the discussion from (3.53) to (3.73), and deduce that the jumps of $(\operatorname{div} u, u_t, \theta, q)$ on Σ_5 and Σ_6 have the same behaviour as in (3.70)-(3.73), and on Σ_k for $k = 7, 8$,

$$\left\{ \begin{aligned} &|[\operatorname{div} u]_{\Sigma_k(t)}| \leq \frac{1}{2\alpha} |\alpha [\operatorname{div} u^0]_\sigma \mp \nabla \Phi^0 \cdot [u^1]_\sigma + \mathcal{O}(\tau)| \\ &\quad \cdot e^{-\frac{\beta\delta}{2\kappa\gamma} t + \left(\frac{1}{2} \sum_{j,l=1}^3 \|\partial_l f_j\|_{L^\infty} \|\partial_{x_j} \Phi^0 \partial_{x_l} \Phi^0\|_{L^\infty} + \mathcal{O}(\tau) \right) t \mp \frac{\alpha}{2} \int_0^t \Delta \Phi^0(x(s; 0, x^0)) ds} \end{aligned} \right. \quad (4.15)$$

$$\left\{ \begin{aligned} &|[\operatorname{div} u]_{\Sigma_k(t)}| \geq \frac{1}{2\alpha} |\alpha [\operatorname{div} u^0]_\sigma \mp \nabla \Phi^0 \cdot [u^1]_\sigma + \mathcal{O}(\tau)| \\ &\quad \cdot e^{-\frac{\beta\delta}{2\kappa\gamma} t \mp \frac{\alpha}{2} \int_0^t (\Delta \Phi^0(x(s; 0, x^0)) + \mathcal{O}(\tau)) ds} \left(2 - e^{\left(\frac{1}{2} \sum_{j,l=1}^3 \|\partial_l f_j\|_{L^\infty} \|\partial_{x_j} \Phi^0 \partial_{x_l} \Phi^0\|_{L^\infty} + \mathcal{O}(\tau) \right) t} \right) \end{aligned} \right.$$

$$\left\{ \begin{aligned} &|[u_t]_{\Sigma_k(t)}| \leq \frac{1}{2} |\nabla \Phi^0 \cdot [u^1]_\sigma \mp \alpha [\operatorname{div} u^0]_\sigma + \mathcal{O}(\tau)| \\ &\quad \cdot e^{-\frac{\beta\delta}{2\kappa\gamma} t + \left(\frac{1}{2} \sum_{j,l=1}^3 \|\partial_l f_j\|_{L^\infty} \|\partial_{x_j} \Phi^0 \partial_{x_l} \Phi^0\|_{L^\infty} + \mathcal{O}(\tau) \right) t \mp \frac{\alpha}{2} \int_0^t \Delta \Phi^0(x(s; 0, x^0)) ds} \end{aligned} \right. \quad (4.16)$$

$$\left\{ \begin{aligned} &|[u_t]_{\Sigma_k(t)}| \geq \frac{1}{2} |\nabla \Phi^0 \cdot [u^1]_\sigma \mp \alpha [\operatorname{div} u^0]_\sigma + \mathcal{O}(\tau)| \\ &\quad \cdot e^{-\frac{\beta\delta}{2\kappa\gamma} t \mp \frac{\alpha}{2} \int_0^t (\Delta \Phi^0(x(s; 0, x^0)) + \mathcal{O}(\tau)) ds} \left(2 - e^{\frac{1}{2} \sum_{j,l=1}^3 \|\partial_l f_j\|_{L^\infty} \|\partial_{x_j} \Phi^0 \partial_{x_l} \Phi^0\|_{L^\infty} + \mathcal{O}(\tau) t} \right) \end{aligned} \right.$$

$$|[\theta]_{\Sigma_k(t)}| \leq \mathcal{O}(\tau) e^{-\frac{\beta\delta}{2\kappa\gamma} t + \left(\frac{1}{2} \sum_{j,l=1}^3 \|\partial_l f_j\|_{L^\infty} \|\partial_{x_j} \Phi^0 \partial_{x_l} \Phi^0\|_{L^\infty} + \mathcal{O}(\tau) \right) t \mp \frac{\alpha}{2} \int_0^t \Delta \Phi^0(x(s; 0, x^0)) ds} \quad (4.17)$$

$$\left\{ \begin{aligned} &|[q]_{\Sigma_k(t)}| \leq \frac{\delta}{2\gamma} |(\nabla \Phi^0 \cdot [u^1]_\sigma \mp \alpha [\operatorname{div} u^0]_\sigma) \nabla \Phi^0 + \mathcal{O}(\tau)| \\ &\quad \cdot e^{-\frac{\beta\delta}{2\kappa\gamma} t + \left(\frac{1}{2} \sum_{j,l=1}^3 \|\partial_l f_j\|_{L^\infty} \|\partial_{x_j} \Phi^0 \partial_{x_l} \Phi^0\|_{L^\infty} + \mathcal{O}(\tau) \right) t \mp \frac{\alpha}{2} \int_0^t \Delta \Phi^0(x(s; 0, x^0)) ds} \end{aligned} \right. \quad (4.18)$$

$$\left\{ \begin{aligned} &|[q]_{\Sigma_k(t)}| \geq \frac{\delta}{2\gamma} |(\nabla \Phi^0 \cdot [u^1]_\sigma \mp \alpha [\operatorname{div} u^0]_\sigma) \nabla \Phi^0 + \mathcal{O}(\tau)| \\ &\quad \cdot e^{-\frac{\beta\delta}{2\kappa\gamma} t \mp \frac{\alpha}{2} \int_0^t (\Delta \Phi^0(x(s; 0, x^0)) + \mathcal{O}(\tau)) ds} \left(2 - e^{\frac{1}{2} \sum_{j,l=1}^3 \|\partial_l f_j\|_{L^\infty} \|\partial_{x_j} \Phi^0 \partial_{x_l} \Phi^0\|_{L^\infty} + \mathcal{O}(\tau) t} \right). \end{aligned} \right.$$

Therefore, similar to Theorem 4.1 of [11], we conclude

Theorem 4.1: For the problem (4.1)-(4.5), in addition to the assumptions of Theorem 3.3, we also suppose that the nonlinear functions f and g are smooth and globally Lipschitz in their arguments. Then, we have

(1) The jumps of $\operatorname{div}u, \partial_t u, \theta, q$ on Σ_5 and Σ_6 decay exponentially both when $\tau \rightarrow 0$ for a fixed $t > 0$ and when $t \rightarrow +\infty$ for a fixed $\tau > 0$.

(2) The jumps of $\operatorname{div}u, \partial_t u, q$ on Σ_7 (Σ_8 resp.) are propagated, and when $t \rightarrow +\infty$ they will decay exponentially as soon as

$$\frac{\beta\delta}{\kappa\gamma} + \alpha\Delta\Phi^0 - \sum_{j,l=1}^3 \|\partial_l f_j\|_{L^\infty} \|\partial_{x_j}\Phi^0 \partial_{x_l}\Phi^0\|_{L^\infty}$$

$$\left(\frac{\beta\delta}{\kappa\gamma} - \alpha\Delta\Phi^0 - \sum_{j,l=1}^3 \|\partial_l f_j\|_{L^\infty} \|\partial_{x_j}\Phi^0 \partial_{x_l}\Phi^0\|_{L^\infty} \quad \text{resp.}\right)$$

is positive, more rapidly for smaller heat conductive coefficient $\kappa\gamma$, while the jump of the temperature θ on Σ_7 and Σ_8 vanishes of order $O(\tau)$ when $\tau \rightarrow 0$.

Remark 4.1: It is easy to carry out the above discussion for the case that the nonlinear functions f and g depending on $\operatorname{div}u$ and q as well in the problem (4.1)-(4.5).

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References

- [1] Beals, M.: *Propagation and Interaction of Singularities in Nonlinear Hyperbolic Problems*. Birkhäuser, Boston 1989.
- [2] Chandrasekharaiah, D. S.: Hyperbolic thermoelasticity: a review of recent literature. *Appl. Mech. Rev.* **51** (1998), 705–729.
- [3] Chen, S. and Wang, Y. G.: Propagation of singularities of solutions to hyperbolic-parabolic coupled systems. *Math. Nachr.*, 242(2002), 46-60.
- [4] Hoff, D.: Dynamics of singularity surfaces for compressible, viscous flows in two space dimensions. *Commun. Pure Appl. Math.* **55** (2002), 1365–1407.
- [5] Gilbarg, D, and Trudinger, N.S.: *Elliptic Partial Differential Equations of Second Order*, Grundlehren math. Wiss. **224**, Springer- Verlag, Berlin (1983).
- [6] Jiang, S. and Racke, R.: *Evolution Equations in Thermoelasticity*, π Monographs Surveys Pure Appl. Math. 112, Chapman & Hall/CRC, Boca Raton (2000).
- [7] Métivier, G.: Propagation, interaction and reflection of discontinuities progressing waves for semi-linear hyperbolic systems. *Amer. J. Math.* **111** (1989), 239–287.

- [8] Racke, R.: Thermoelasticity with second sound — exponential stability in linear and non-linear 1-d. *Math. Meth. Appl. Sci.* **25** (2002), 409–441.
- [9] Racke, R.: Asymptotic behavior of solutions in linear 2- or 3-d thermoelasticity with second sound. *Quart. Appl. Math.* **61** (2003), 315–328.
- [10] Racke, R. and Wang, Y. G.: Propagation of singularities in one-dimensional thermoelasticity. *J. Math. Anal. Appl.* **223** (1998), 216–247.
- [11] Racke, R. and Wang, Y. G.: Asymptotic behavior of discontinuous solutions to thermoelastic systems with second sound. *J. Anal. Appl.* **24** (2005), 117–135.
- [12] Rauch, J. and Reed, M.: Discontinuous progressing waves for semilinear systems. *Commun. PDE* **10** (1985), 1033–1075.
- [13] Wang, Y.G.: Microlocal analysis in nonlinear thermoelasticity. *Nonlinear Anal.* **54** (2003), 683–705.
- [14] Wang, Y.G.: A new approach to study hyperbolic-parabolic coupled systems. In “*Evolution Equations*” (R. Picard, M. Reissig & W. Zajaczkowski eds.), *Banach Center Publications*, 60(2003), 227–236.